

# Measuring the brightness temperature of the sun using a DSTV dish

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# 1. Introduction

## 1.1 Theory

### 1.1.1 Radio waves

***Radio waves -> electromagnetic signals having wavelengths greater than a centimetre (see Fig 1. below).***

E.g.

- Commercial **FM radio** operates in the frequency range from **88 to 108 MHz**, ( $\lambda \sim 3$  metres).
- **Cellphones** operate at a frequency  $\nu$  of about **900 MHz**, ( $\lambda \sim 33$  cm).

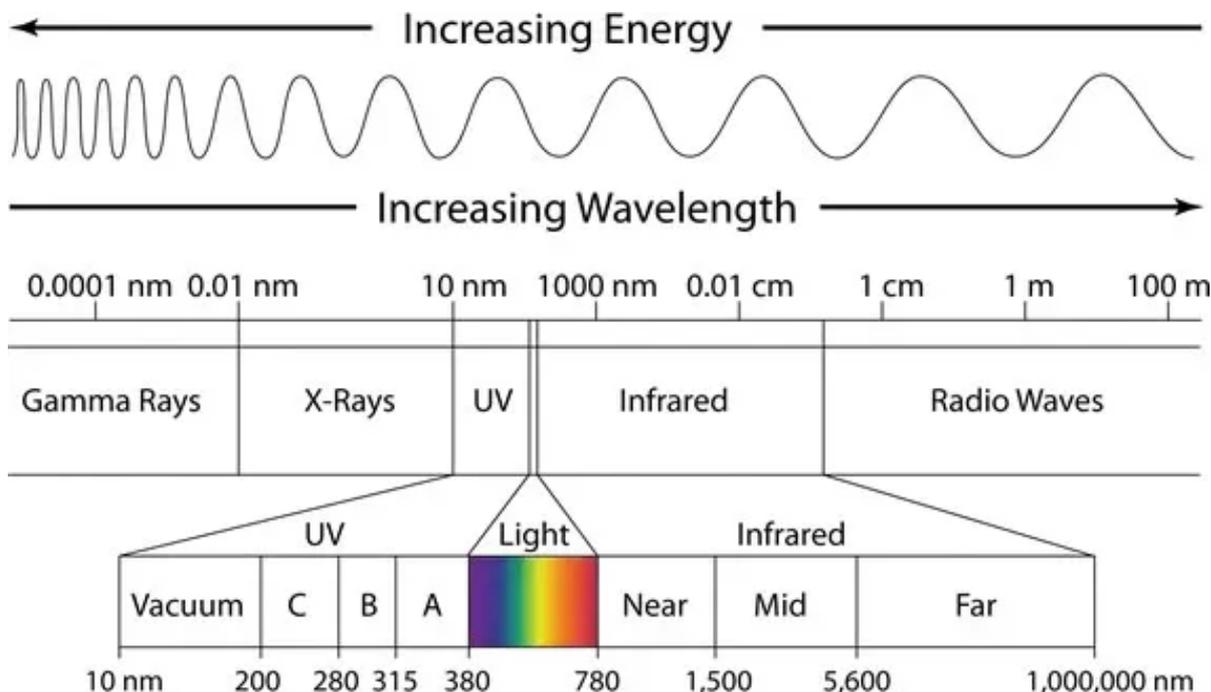


Fig 1. credit: <https://i0.wp.com/apniphysics.com/wp-content/uploads/2023/10/electromagnetic-spectrum-wavelengths.jpg?ssl=1>  
(<https://i0.wp.com/apniphysics.com/wp-content/uploads/2023/10/electromagnetic-spectrum-wavelengths.jpg?ssl=1>)

**The microwave band is the short wavelength part of the radio band and covers 1 to 30 cm wavelength.**

- **Microwave ovens** operate at **2.4 GHz** ( $\lambda \sim 12$  cm).
- **DSTV satellites** transmit at **12 GHz** ( $\lambda \sim 2.5$  cm).

**Why are cellphones are not allowed at a radio observatory ?**

With the advent of technology comes great consequences

## Frequency Bands by Generation

### 2G (GSM):

Primarily uses frequencies around 850 MHz, 900 MHz, 1800 MHz, and 1900 MHz. [↗](#)

### 3G (UMTS):

Operates in a wider range, including 850 MHz, 900 MHz, 1700 MHz, 1900 MHz, and 2100 MHz. [↗](#)

### 4G (LTE):

Employs a broad spectrum, from lower bands (like 600 MHz and 700 MHz) to mid-bands up to around 2.6 GHz and even some higher frequencies. [↗](#)

### 5G:

Utilizes three primary types of spectrum:

- **Low-band:** Similar to 4G frequencies (below 1 GHz) for broad coverage.
- **Mid-band:** Offers a balance of speed and coverage, often between 1 GHz and 6 GHz.
- **High-band (mmWave):** Uses millimeter-wave frequencies (24 GHz to 100 GHz) for extremely high data rates but with shorter range. [↗](#)

Band	Frequency Range	Bandwidth
<b>Mid Band 1</b>	<b>0.35 – 1.05 GHz</b>	<b>700 MHz</b>
<b>Mid Band 2</b>	<b>0.95 – 1.76 GHz</b>	<b>810 MHz</b>
Mid Band 3	1.65 – 3.05 GHz	1.4 GHz
Mid Band 4	2.80 – 5.18 GHz	2.4 GHz
<b>Mid Band 5a</b>	<b>4.6 – 8.5 GHz</b>	<b>3.9 GHz</b>
<b>Mid Band 5b</b>	<b>8.3 – 15.4 GHz</b>	<b>2 x 2.5 GHz</b>

*65k channels maximum across any band, zoom windows possible with 16k channels*

SKA-mid operating freqs.

## 1.1..2 Radio Telescope Antennas (resolution)

**radio telescope** focuses EM waves with a dish and measures their power with a **radiometer**.

A “classic” radio telescope for use in the microwave band has a circular parabolic reflector with a feed horn at the focus to collect the incoming microwaves and pass them to transistor amplifiers in the receiver.

A DSTV satellite dish also works in this way.

- can be used as a **mini-radio telescope by replacing the DSTV decoder with a radiometer for measuring the signal strength.**

DSTV dishes have a smooth solid surface in order to reflect incoming radio waves with high efficiency.

Satellite dishes working at longer wavelengths e.g. 8cm (3.8 GHz) can get away with a rougher mesh surface and still have acceptable efficiency.

To understand how a **reflector antenna responds to radiation** coming from different angles, consider what happens when a plane wave of wavelength  $\lambda$  arrives at a circular aperture of diameter  $D$  (Fig. 2).

**Constructive and destructive interference produces a circularly symmetric diffraction** pattern, with a central maximum and concentric rings of decreasing strength (Fig. 3).

This same pattern describes the **response of a circular antenna to plane waves** coming from different angles, and it is then called an **antenna beam pattern**.

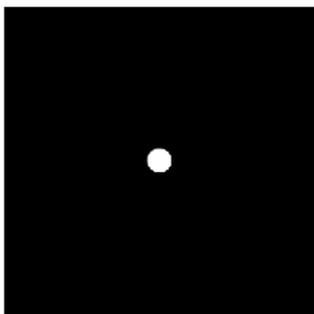


Figure 2: A focusing lens or reflector with a circular aperture.

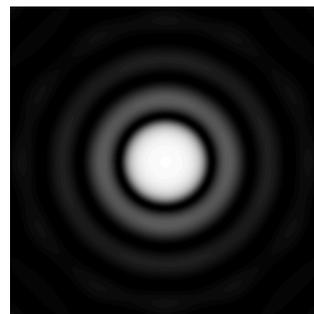


Figure 3: The diffraction pattern produced by a circular focusing lens or reflector.

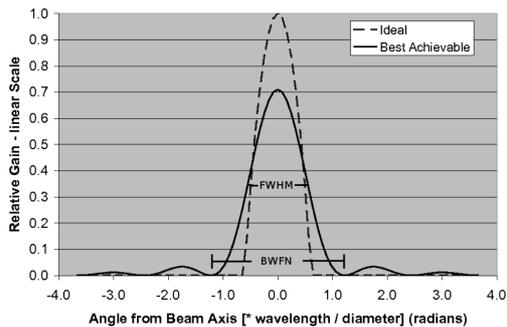


Figure 4: Cross-sections through the beam pattern of an ideal antenna and a practically realizable antenna, shown with a linear vertical scale.

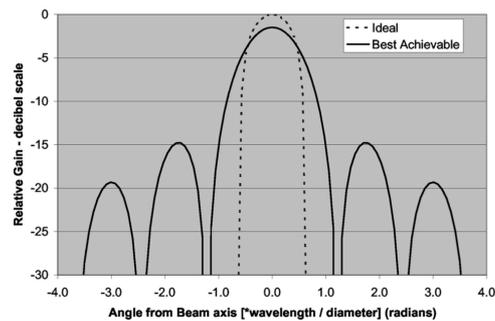


Figure 5: Beam cross-sections of ideal and practically realizable antennas, with a logarithmic vertical scale to show the sidelobe structure.

- An “ideal” antenna would produce a beam that captures 100% of the incoming energy in the main beam and would have no sidelobes.
- => “main beam efficiency”  $\epsilon_m$  of 1.0.
- not possible to actually achieve this, and  $\epsilon_m$  **usually lies between 0.6 and 0.8**.

Figs. 4 and 5 show an ideal and an actual beam pattern in cross-section on linear and logarithmic scales.

The “ideal” pattern has been modelled here with a parabolic shape, while the mathematical form of the real pattern is a  $\text{sinc}^2$  or  $\left(\frac{\sin(x)}{x}\right)^2$  function.

- This describes the diffraction pattern where nothing obstructs the path of the waves, i.e. it has an “**unblocked** aperture”, and is **uniformly illuminated**.
- The **first minimum or null in the pattern occurs at a radius of about  $1.2 \frac{\lambda}{D}$  radians**, so the beamwidth to first nulls is

$$BWFN \sim 2.4 \frac{\lambda}{D} [\text{radians}] \quad (1)$$

**Antenna beam / HPBW:** resolution - angular size of the main lobe of the beam at the points where the power drops to half of the maximum value.

- The beamwidth at the half-power points (HPBW), also called the Full Width at Half Maximum (FWHM), is about half this, as shown in figs. 4 and 5.

$$HPBW = FWHM \sim BWFN/2 = 1.2 \frac{\lambda}{D} [\text{radians}] \quad (2)$$

### 1.1.3 Brightness Temperature and Antenna Temperature

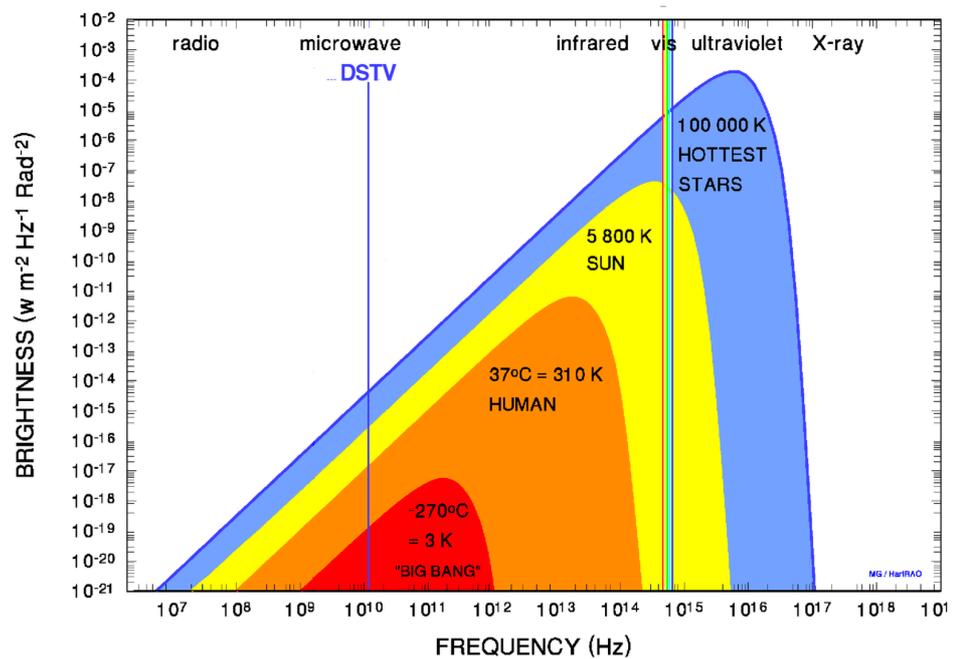
- Brightness temperature: translates “how strong the radiation is” into an equivalent blackbody temperature.
- When we measure radio emission, what we directly detect is flux density (how much power per area per frequency).
- But astronomers often want a way to describe this emission in temperature terms, because thermal sources (like stars, planets, dust) naturally radiate according to temperature.
- Brightness temperature is a way to express the intensity of radiation as if it were coming from a blackbody at that temperature.

Fig. 6 shows the brightness as a function of frequency for several black body radiators modelled as having equal size but different temperatures.

The frequencies of satellite TV transmission and visible light are marked.

Clearly, \*\*hotter objects produce more radiation than cooler ones, and the brightness maximum occurs at a higher frequency. \*\*

The wavelength or frequency at which the intensity peaks is given by the well-known **Wien displacement law**.



*Blackbody radiation from solid objects of the same angular size, at different temperatures*

From Fig. 6 we can see that for all objects with temperatures more than a few degrees above absolute zero, the brightness peak occurs well above the operating range of radio telescopes.

- Hence radio telescopes work in the range where  $h\nu \ll kT$ , so the
- Rayleigh-Jeans law applies and the brightness  $B$  - and hence the power measured by a radio telescope - is proportional to the temperature  $T$  of the emitting source:

$$B = \frac{2kT}{\lambda^2} \quad [W m^{-2} Hz^{-1} sr^{-1}] \quad (3)$$

where  $k$  = Boltzmann's constant =  $1.38 \times 10^{-23} [JK^{-1}]$ .

The apparent temperature of an emitting source at a given frequency is a property of the object emitting the radio waves. This is called its "**brightness temperature**",  $T_B$ .

For some astronomical objects the brightness temperature that is measured using a radio telescope is meaningful as a physical temperature, for example when observing a planet or moon. For other objects it may not be, depending on the mechanism that produced the radio emission.

**By pointing the antenna at objects of known temperature that completely fill the beam, we can calibrate the output signal in units of absolute temperature (Kelvins). So one can think of a radio telescope as a remote-sensing thermometer.**

- What does the telescope measure ?

The "**antenna temperature**"  $T_A$  of a source is the increase in temperature (receiver output) measured when the antenna is pointed at a radio-emitting source. It will be less than brightness temperature of the source if the source does not fill the whole beam of the telescope. Note that "antenna temperature" has nothing to do with the physical temperature of the antenna.

- How do we obtain it ?

**To obtain the brightness temperature  $T_B$  of the emitting source from its measured antenna temperature  $T_A$ , we have to measure the angular size of the source and of the telescope beam.**

**The ratio of the angular size (solid angle) of the source to the angular size of the beam gives the fraction of the beam that is filled by the source. aka. a scaling factor.**

This will be useful to remember during the experiment later on.

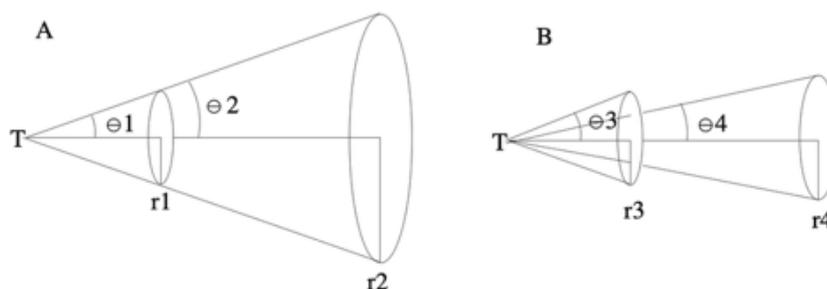


Figure 7: Comparison of physical radius  $r$  and angular radius  $\theta$ . In diagram A the two objects have different physical radii:  $r_2 > r_1$ , but the same angular radii as seen by telescope T:  $\theta_1 = \theta_2$ . This applies to the Moon and Sun as seen from Earth. By contrast, in diagram B,  $r_4 > r_3$ , but  $\theta_4 < \theta_3$ .

### 1.1.4 Angular Sizes and the Sun's Brightness Temperature

The difference between physical diameter and angular diameter is shown in Fig. 7.

For example, the Sun has a physical diameter of 1.4 million km, while the Moon has a diameter of 3500 km.

Yet, as seen from the Earth, the Sun and the Moon appear to be the same size, i.e. they have the same angular diameter.

#### How can this be ?

The Sun is 400 times bigger than the Moon, but it is also 400 times further away.

As the projection of the antenna beam onto the sky is two-dimensional, we shall need to find the angular area that it covers.

Angular area is called a “solid angle” and the units are radians<sup>2</sup>, or steradians (sr).

An object with an angular radius  $\theta$  radians subtends a solid angle.

$$\Omega = 2\pi(1 - \cos\theta) \quad [sr] \quad (4)$$

For small  $\theta$ ,

$$\Omega = \pi\theta^2 \quad [sr] \quad (5)$$

We can use this equation to calculate the solid angle of the Sun as seen from the Earth,  $\Omega_s$ .

The beam solid angle  $\Omega_A$  of the antenna can be obtained from the half-power beamwidth (HPBW) - in units of radians - by assuming the main lobe of the beam has a Gaussian shape:

$$\Omega_A \sim 1.333(HPBW)^2 \quad [sr] \quad (6)$$

For small  $\theta$ , we can assume that the

$$T_B = \frac{\Omega_A T_A}{\Omega_s \epsilon_m} \quad [K] \quad (7)$$

## 1.2. Detecting Radio Emission from Space

When the telescope looks at a radio source in the sky, the receiver output is the sum of the radio waves received from several different sources:

- Behind the radio source whose brightness we want to measure is the **cosmic microwave background (CMB)** coming from every direction in space. This is the radiation emitted as the first atoms formed, 380 000 years after the Big Bang. The black body temperature of the CMB  $T_{cmb}$  has now decreased to **2.7 Kelvins**, as the expansion of the Universe has stretched out the waves by a factor of 1000. The CMB produces a brightness temperature  $T_{Bcmb}$  of  $\sim 2.7$  K at 1.4 or 4 GHz, reducing to 2.5 K at 12 GHz.
- The emission from the radio source we want to measure, which produces the antenna temperature  $T_A$ .
- Radiation from the dry atmosphere  $T_{at}$  and
- from the water vapour in the atmosphere  $T_{wv}$ .

The dry air adds about 1K, and at 12 GHz water vapour adds 1 - 2 K, depending on the humidity.

- The radiation the feed receives through the antenna sidelobes from the (warm - 290 K) ground or nearby buildings beyond the edge of the antenna, of brightness temperature  $T_g$ . With the antenna pointing straight up at zenith this could add 5 - 15 K; 10 K is a reasonable number to use. It increases when pointing close to the horizon.

- The amplifiers in the receiver generate their own electronic noise and so produce a receiver noise temperature  $T_R$ .

The sum of these parts is called the “system temperature”  $T_{sys}$ . Summing from the most distant noise contributor to the nearest we have:

$$T_{sys} = T_{Bcmb} + T_A + T_{at} + T_{wv} + T_g + T_R \quad [K] \quad (8)$$

## 2. The Experiments

**In this experiment we are going to determine the brightness temperature of the sun by hand to give you hands on experience on how this can be done with very little equipment.**

The aim is to estimate

$$T_B = \frac{\Omega_A T_A}{\Omega_s \epsilon_m} \quad [K] \quad (7)$$

To do this, we need to find:

1. The antenna beam width  $\Omega_A$
2. The angular diameter of the sun  $\Omega_s$
3. The antenna temperature  $T_A$
4. The temperture brightness  $T_B$

We determine steps 1 - 3 in and use those results to calculate 4.

1. We calibrate our equipment to get a conversion factor (c) [DSTV dish Hot/Cold load experiment]
2. Measure the temperature of the sun. [DSTV dish Hot/Cold load experiment]
3. Measure the angular diameter of the sun [Pinhole experiment]

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You are now going to use the principles discussed in the previous section to calculate the brightness temperature of the Sun. But before you start, we need to setup our experiment. Here's what we are going to need:

### 2.1 Equipment

### 2.1.1. A standard small Ku-band (12GHz) satellite dish.

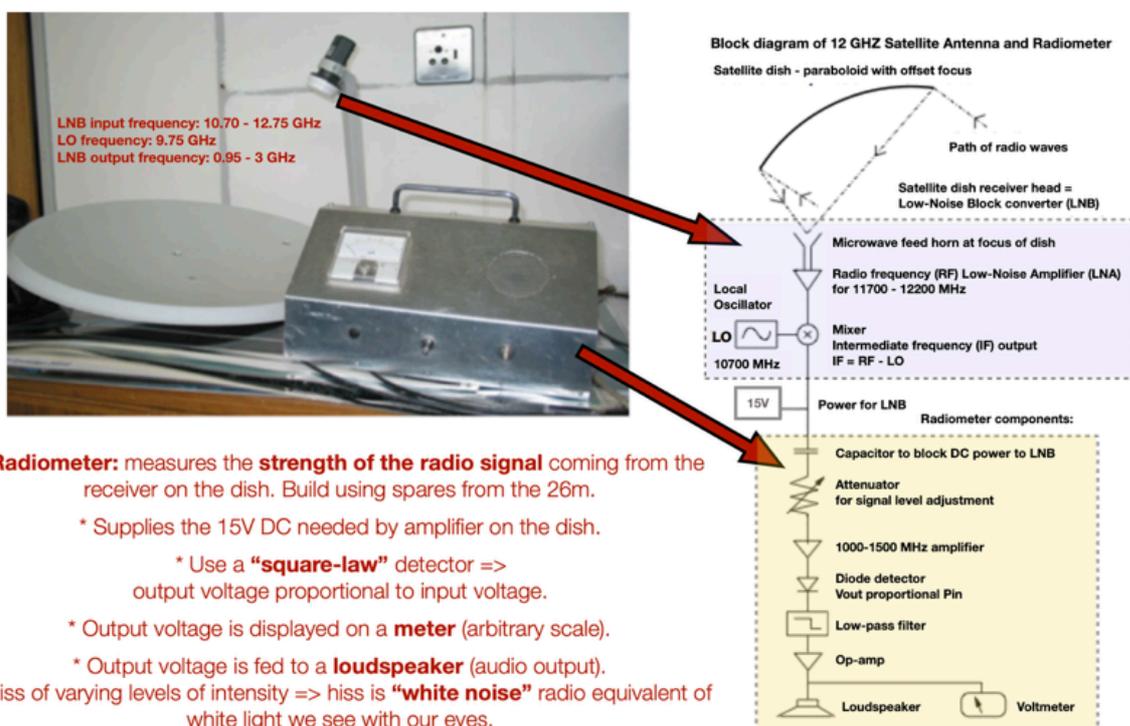
The main parts of a simple radio telescope comprising a satellite dish and radiometer are shown in the figure below.

The dish output is fed to a radiometer. A **“radiometer”** means a **“radio-meter”** – i.e. a **device that measures the strength of the radio signal coming from the receiver on the dish.**

Since the incoming radio waves from natural emitters are weak and noise-like.

If the output of the detector is connected to a loudspeaker, the signal sounds like a hiss, as one hears if a radio is tuned off-station.

The internal noise produced in the amplifiers is generally larger than the signal from natural radio sources.



### 2.1.2. Apparatus for Measuring the Diameter of the Sun

You will need:

- two pieces of card or paper,
- a pencil,
- a sharp knife,
- a ruler and
- a 2 metre tape measure.

### 2.1.3. Apparatus for Measuring the Antenna Beamwidth

You will need:

- pen and
- notepad or paper.
- (Optionally): Apparatus required is a tripod or mount on which the satellite dish can be locked in position, a timer (e.g. watch on which seconds are displayed). We wont be using this for our experiment.

```
In [1]: 1 # 20*254
        2 1.2*(2.5/50)*57.3
```

```
Out [1]: 3.4379999999999997
```

## 2.2 Experimental procedure

The temperature brightness  $T_B$  is described in eq. (7) as

$$T_B = \frac{\Omega_A T_A}{\Omega_s \epsilon_m} \quad [K]$$

where  $\Omega_A$  is the beam solid angle,  $\Omega_s$  is the solid angle of the Sun as seen from Earth,  $\epsilon_m$  is the main beam efficiency and  $T_A$  is the emission from the radio source we want to measure. In order to measure  $T_A$  we need to establish a conversion factor or "Calibrate our telescope" so that we can be able to express our detected radio emission it in terms of temperature in Kelvins [K].

### 2.2.1 Calibrating the Radio Telescope

When a radio telescope detects radio emission from a celestial object, the dish produces an output voltage proportional to the object's radio power (a "square-law" detector) plus its own internal receiver noise.

- This voltage is displayed on a meter with an arbitrary scale.
- To get the temperature brightness of the object, we need a way to establish a scale of Kelvins per radiometer output unit. To establish a scale of Kelvins per radiometer output unit, we use a technique that leverages the sky and ground as reference points.

#### *Determining the Scaling Factor*

We can determine this scaling factor by using the sky as a "cold load" and the ground as a "hot load."

- **Cold Load (Sky):** By pointing the telescope at the zenith, we effectively observe a "cold load" with a temperature  $T_{sky} \sim 10$  K. This temperature accounts for the Cosmic Microwave Background, atmospheric contributions, and ground radiation.
- **Hot Load (Ground):** Aiming the telescope at the ground provides a "hot load" with a temperature  $T_g \sim 300$ K on a warm day.

Let:

- $V_1$ : Meter reading for the sky
- $V_2$ : Meter reading for the ground
- $T_R$ : Receiver noise temperature
- $T_{sky}$ : Effective sky temperature
- $T_g$ : Ground temperature
- $c$ : Constant of proportionality

if  $V_1$  and  $V_2$  are the meter readings for the sky and ground, respectively. The relationship between temperature and voltage can be expressed as:

$$T_R + T_{sky} = cV_1 \quad [K] \quad (9)$$

where  $T_{sky} = T_{Bcmb} + T_{at} + T_{wv} + T_g \sim 10$ K.

$$T_R + T_g = cV_2 \quad [K] \quad (10)$$

By solving these equations simultaneously, we can determine the value of  $c$  and establish a calibrated temperature scale for the radiometer.

### **HartRAO Satellite dish Radiometer**

On the HartRAO satellite dish radiometer, typical meter readings are:

10: Aimed at the sky ( $V_1$ )

30: Aimed at the ground ( $V_2$ )

24: Aimed at the Sun ( $V_3$ ), you will need this for the next calculation

By using these readings and the known temperatures of the sky and ground, we can calculate the scaling factor  $c$  and subsequently determine the temperature of the Sun (or any other celestial object) based on its corresponding meter reading.

**>>> Solution for Determining the Scaling Factor**

Step 1: Establish the equations:

As given in the experiment notes, we have two equations:

$$T_R + T_{sky} = cV_1 \quad [K] \quad (9)$$

$$T_R + T_{ground} = cV_2 \quad [K] \quad (10)$$

Step 2: Solve for the scaling factor,  $c$ :

To find the scaling factor, we can subtract equation 9 from equation 10:

$$(T_R + T_g) - (T_R + T_{sky}) = cV_2 - cV_1$$

This simplifies to:

$$(T_g - T_{sky}) = c(V_2 - V_1)$$

Therefore, the scaling factor,  $c$ , is:

$$c = \frac{T_g - T_{sky}}{V_2 - V_1}$$

Step 3: Calculate the scaling factor using given values:

$$c = \frac{300K - 10K}{30 - 10} = \frac{290K}{20} = 14.5 \quad K/unit$$

In [2]:

```

1 Tg=300
2 Tsky=10
3 V1=10
4 V2=30
5
6 c= (Tg-Tsky)/(V2-V1)
7
8 f'The scaling factor is {c} kelvins per arbitrary unit'
```

Out[2]: 'The scaling factor is 14.5 kelvins per arbitrary unit'

Now that we know  $c$  we can determine the noise contribution from  $T_R$

substituting for  $c$  in equation (9)

$$14.5(10) = T_R + 10$$

$$T_R = 145 - 10$$

$$T_R = 135 \quad [K]$$

```
In [3]: 1 TR = int((c*V1)-Tsky)
        2
        3 f'The noise temperature of the reciever is {TR} K'
```

```
Out[3]: 'The noise temperature of the reciever is 135 K'
```

## 2.2.2 Measure the Antenna Temperature from the Sun

We need to aim the telescope at the Sun. With an offset-fed paraboloid this can be a little tricky. Hold the dish horizontal. then turn the dish horizontally so the shadow of the feed arm falls across the centre of the dish, demonstrated here by Mike in the image below.



Looking past the feed to see the reflection of the Sun off the dish surface

Then rotate the dish in elevation so the shadow of the feed arm on the dish gets shorter. The Sun comes into the beam just after the shadow leaves the edge of the dish. Adjust the direction gently up and down and sideways to maximise the signal from the Sun; the signal should roughly double. Once the telescope is accurately pointed at the Sun, record the new meter reading,  $V_3$ . This reading corresponds to the combined temperature of the receiver, sky, and Sun, and is expressed as:

$$cV_3 = T_R + T_{sky} + T_{Asun} \quad (11)$$

Calculating the Sun's Antenna Temperature:

Since the scaling factor,  $c$ , and the receiver temperature,  $T_R$ , are known from the calibration process, and we have a reasonable estimate of the sky temperature,  $T_{sky}$ , we can use equation 11 to directly calculate the antenna temperature of the Sun,  $T_{Asun}$ .

$$T_{Asun} = cV_3 - T_R - T_{sky}$$

**>>> Solution for Measuring the Antenna Temperature from the Sun**

$$T_{Asun} = (14.5K/unit)(24units) - 135K - 10K \quad (12)$$

$$T_{Asun} = 348 - 145K$$

$$T_{Asun} = 203 K$$

```
In [4]: 1 V3=24
        2
        3 TAsun = int(c*V3) - TR - Tsky
        4
        5 f'The antenna temperature of the sun is {TAsun} K'
```

Out[4]: 'The antenna temperature of the sun is 203 K'

But the Sun is much hotter than this! To understand why this is so, and to work out the real temperature of the Sun, we must consider the telescope's field of view. We have to figure out how much of the sky the dish “sees” (or more technically, the beam solid angle) at a given time.

**2.2.3 Measure the Angular Diameter of the telescope beam**

## 1. Angular diameter or angular resolution of The telescope beam

A radio telescope, like the satellite dish you're using, is sensitive to radio waves originating from a specific region of the sky. This region is known as the antenna's "beam."

### Understanding Beamwidth

What is the diameter of this beam ?

As the radio source moves away from the middle of the beam, the angle of the waves hitting the dish changes.

When all the waves coming from each part of the dish are in phase, we get the strongest signal, when the emitter is in the centre of the beam. Moving away from the centre, the waves from different parts of the dish start to become out of phase with each other, causing destructive interference.

The telescope sensitivity falls to a minimum when the incoming wave is tilted at an angle so that there is a phase difference of approximately one wavelength across the diameter of the dish.

This minimum is called the "first null" in the antenna beam pattern. We normally use the points where the response of the antenna has fallen to half that in the centre of the beam defined as the "half-power beamwidth" ( $\theta_{HPBW}$ ) or angular resolution.

$\theta_{HPBW} \sim \text{wavelength} / \text{diameter of the dish, in units of radians (1 radian = 57.3 degrees)}$ .

For a circular aperture, like our dish, the angular resolution,  $\theta_{HPBW}$ , can be estimated using the Rayleigh criterion given in equation (2).

- Calculate the HPBW of the telescope beam.

**For comparison, also calculate the resolving power of your eye (2 mm pupil) at a wavelength of 500 nm).**

### HartRAO Satellite dish parameters

dish\_diameter = 50 # cm, satellite dish diameter

wavLen = 2.5 # cm, observing wavelength

### >>> Solution for measuring the Angular diameter or angular resolution of The telescope beam

1. Calculation of the  $\theta_{HPBW}$

$$\theta_{HPBW} = (1.2)(obs\_freq)/dish\_diameter$$

$$\theta_{HPBW} = (1.2)(2.5)/50$$

$$\theta_{HPBW} = 0.060 \quad [radians]$$

$$= 3.4 \text{ degrees (approx)}$$

```
In [5]: 1 import numpy as np
        2
        3 dish_diameter = 50
        4 wavLen = 2.5
        5 theta_hpwb_rad = (1.2*wavLen)/dish_diameter
        6 theta_hpwb = np.rad2deg(theta_hpwb_rad) # convert radians to de
        7
        8 f'The hpbw or angular diameter of the antenna beam is {theta_hp
```

```
Out[5]: 'The hpbw or angular diameter of the antenna beam is 0.06 radians
or ~3.44 degrees'
```

This is the beam of the telescope. We now need to figure out what is the beam solid angle  $\Omega_A$  or the angle subtended in the sky by the telescope beam.

This is a crucial parameter in radio astronomy as it determines the amount of celestial radiation that the telescope can collect or more technically a measure of the telescope's field of view on the sky.

A larger beam solid angle means the telescope collects radiation from a larger area of the sky, while a smaller beam solid angle provides a more focused view.

Calculate the antenna beam solid angle  $\Omega_A$

### >>> Solution for measuring the solid angle or area of The telescope beam

1. Calculation of the solid angle  $\Omega_A$

The beam area (solid angle) is agiven in eqn. (6)

$$\Omega_A = 1.133(\theta_{hpwb})^2 \quad [sr]$$

$$\Omega_A = 1.133(3.44)^2 \quad [sr]$$

$$= 13.39$$

```
In [6]: 1 import math
        2
        3 omega_a_deg = 1.133 * (theta_hpwb)**2
        4
        5 f'The solid angle or area of the antenna beam is {omega_a_deg:.
```

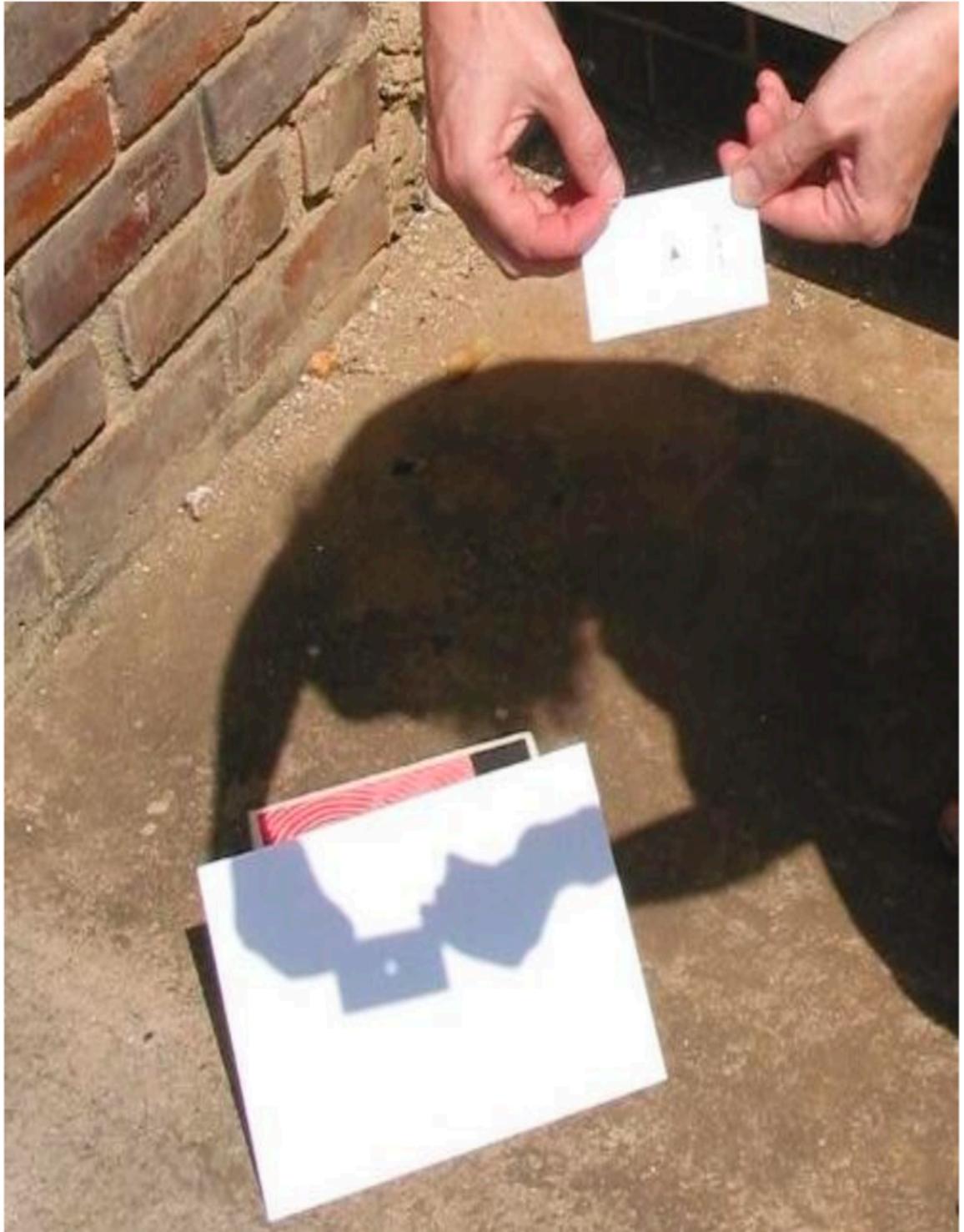
```
Out[6]: 'The solid angle or area of the antenna beam is 13.39 square degrees'
```

## 2.2.4 Measure the Angular diameter of the Sun

To measure the angular diameter of the Sun, we will use the pinhole projection experiment. Assuming that our target, the Sun, is much smaller or can fit within the beam of our antenna, we now want to find out what its angular diameter is.

Put simply, we need to know how many times the Sun would fit into the beam of the telescope.

### 1. Estimating the temperature of the Sun using the beam size



Projecting the Sun through a pinhole.

Use pinhole projection to measure the angular diameter of the Sun.

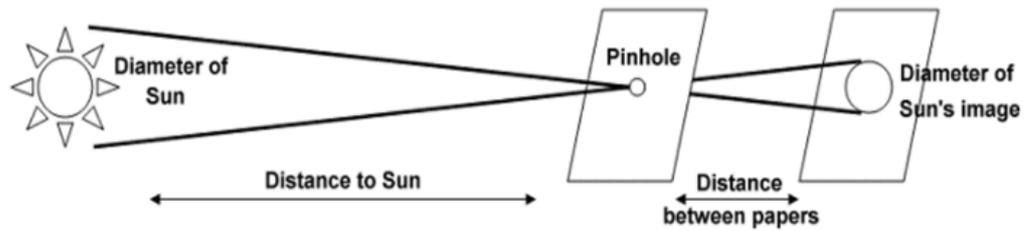
On one piece of card use a ruler and pencil to mark three equilateral triangles several cm apart, with sides of about two, three and four millimetres.

Cut out the triangles using the sharp knife and ruler.

The card with the holes is used to project images of the Sun onto the second card (Fig above and below).

The two cards need to be separated by a distance  $D$ , measured by the tape measure.

**If the dish can be pointed accurately using a tripod, the beamwidth could actually be measured, by scanning across the Sun or letting it drift through the beam. This is one of the methods used to observe radio sources by the HartRAO 26m telescope.**



1. Stand up and hold the string or tape measure in one hand. Stretch your arm straight up, let the string touch the ground. Tie a knot in the string where it touches the ground. Use the ruler to measure the length of the string up to the knot.
2. Record the length: \_\_\_\_\_mm
3. Make a hole in the centre of one of the pieces of paper. The hole should be between 2 and 4mm in diameter.
4. Stick the second piece of paper to the outside cover of a hardcover book. This will allow the paper to stand on it's own on the ground.
5. Set the distance between the two pieces of paper as the length of string to the knot.
6. Let the Sun shine through the pinhole onto the second sheet.
7. Draw a circle around the Sun's image on the second sheet. Use the ruler to measure the diameter of the Sun's image: \_\_\_\_\_mm.



8. Divide the answer from number 2 by the answer from number 7. \_\_\_\_\_

This gives us the number of Sun images that when placed end to end, would fit in the distance between the pinhole paper and the graph paper.

9. Using the answer from number 8, we can say that the distance from the Earth to the Sun is \_\_\_\_\_ times the diameter of the Sun. We are using the principle of simple proportion to do this – see the diagram on the next page.

**HartRAO Satellite dish parameters**

D = 2000: mm, length of distance between the 2 papers or cards (from 2, above) measured with a tape measure

d = 20: mm, The linear diameter of the projected Sun measured with the ruler.

**>>> Solution for measuring the Angular diameter or angular resolution of The telescope beam**

The angular diameter of the Sun is  $d/D$  radians:

$$\theta_{sun} = \frac{d}{D} \quad [radians]$$

$$\theta_{sun} = \frac{20}{200} \quad [radians]$$

$$\theta_{sun} = 0.01 \quad [radians]$$

or

$$\theta_{sun} = 0.57^\circ$$

```
In [7]: 1 d=20
        2 D=2000
        3 theta_sun = d/D
        4 theta_sun_deg = np.rad2deg(theta_sun)
        5 f'The angular diameter of the Sun = {theta_sun} [radians] or {t
```

```
Out[7]: 'The angular diameter of the Sun = 0.01 [radians] or 0.57 degrees'
```

**2. Estimating the solar radius of the Sun using the angular diameter**

The solar radius is given as

$$\theta_s = \frac{d}{2D} \quad [radians]$$

$$\theta_s = \frac{20}{2*2000} \quad [radians]$$

```
In [8]: 1 d=20
        2 D=2000
        3 theta_s = d/(2*D)
        4 theta_s_deg = np.rad2deg(theta_s)
        5 f'The solar radius of the Sun = {theta_s} [radians] or {theta_s
```

```
Out[8]: 'The solar radius of the Sun = 0.005 [radians] or 0.29 degrees'
```

### 3. Estimating the angle subtended by the Sun using the solar radius

The angle subtended by the sun or are of the sun is given by equation 5

$$\Omega_s = \pi\theta^2 \quad [sr]$$

$$\Omega_s = \pi(\theta_{s\_deg})^2$$

$$\Omega_s = \pi(0.29)^2$$

$$\Omega_s = 0.26 \quad [degrees \ squared]$$

**Note** If you perform this calculation using the radian results from the previous calculation, your units will be radians squared for the answer. Since we used degrees, our answer is stated in squared degrees.

```
In [9]: 1
        2 omega_s_deg = math.pi*(theta_s_deg)**2
        3
        4 f'The angle subtended by the Sun = {omega_s_deg:.2} degrees squ
```

```
Out[9]: 'The angle subtended by the Sun = 0.26 degrees squared'
```

### 4. Find ratio of the areas of the beams

```
In [10]: 1
         2
         3 area_ratios = int(omega_a_deg/omega_s_deg)
         4
         5 f' and so put simply we could fit {area_ratios} Suns into the b
         6
```

```
Out[10]: ' and so put simply we could fit 51 Suns into the beam of the dis
h.'
```

## 2.2.5 Calculate the Brightness Temperature of the Sun

To calculate the brightness temperature of the Sun we now have the Sun's solid angle  $\Omega_s$  from Eqn. 5, the beam solid angle  $\Omega_A$  from Eqn. 6, and the antenna temperature measured from the Sun  $T_A$ , from Eqn. 11.

However, we need to correct the antenna temperature to reflect the area ratio calculated above

```
In [11]: 1 TAsun_corrected=area_ratios*TAsun
          2 f'Hence the temperature of the sun is about {TAsun_corrected} K'
```

```
Out[11]: 'Hence the temperature of the sun is about 10353 K'
```

This does not allow for the efficiency of the dish. Experiments indicate that a reasonable value for  $\eta_m$  for a DSTV dish at 12 GHz is about 0.75.

For mesh surface satellite dishes working at 3.8 GHz, a lower value is more likely as the surface shape is less accurate, and  $\eta_m$  is about 0.5.

For a mesh dish with a 1.4GHz can feed,  $\eta_m$  is about 0.5.

These values are based on measurements with domestic satellite dishes.

Correcting for this, the Sun's brightness temperature is about

```
In [12]: 1 Tb=int(TAsun_corrected/0.75)
          2 f'So the brightness temperature of the sun can be estimated as'
```

```
Out[12]: 'So the brightness temperature of the sun can be estimated as 13804 Kelvins at 12 GHz'
```

Estimate the uncertainty in your result in the usual way, by propagating estimates of the error in each of your measurements. Assume an uncertainty in  $\epsilon_m$  of 10%.

This temperature will vary during the 11 year solar cycle, but it reasonably matches the value given in standard references for this observing wavelength. However, the  $T_B$  is about three times higher than the value of 5800 K measured at visible wavelengths.

How does your result for the Sun's brightness temperature compare to the temperature usually quoted for the Sun's photosphere (light emitting surface)? What do you think this implies?

In essence, we cannot treat the Sun as a perfect blackbody radiator – some other mechanism must be increasing the emission from the Sun at radio wavelengths. What could this be?

```
In [ ]: 1
```